

Lecture 03 Infinite Square Well



Wave equation for moving particles

... in one of the next colloquia [early in 1926], Schrödinger gave a beautifully clear account of how de Broglie associated a wave with a particle and how he [i.e., de Broglie] could obtain the quantization rules...by demanding that an integer number of waves should be fitted along a stationary orbit. When he had finished Debye² casually remarked that he thought this way of talking was rather childish...[that to] deal properly with waves, one had to have a wave equation.













 $-\frac{\hbar^2}{2m}\frac{\partial^2\Psi(x,t)}{\partial x^2}+V(x,t)$



$$(x,t)\Psi(x,t) = i\hbar \frac{\partial \Psi(x,t)}{\partial t}$$



Separation of the Time and Space Dependencies of $\psi(x, t)$

$$\frac{-\hbar^2}{2m}\frac{d^2\psi(x)}{dx^2} + V(x)\psi(x)$$

$$\Psi^*(x,t)\Psi(x,t) = \psi^*(x)e^{+iEt/\hbar}\psi(x)e^{-iEt/\hbar} = \psi^*(x)\psi(x)$$

- 1. $\psi(x)$ must exist and satisfy the Schrödinger equation. 2. $\psi(x)$ and $d\psi/dx$ must be continuous.
- 3. $\psi(x)$ and $d\psi/dx$ must be finite.
- 4. $\psi(x)$ and $d\psi/dx$ must be single valued.
- 5. $\psi(x) \rightarrow 0$ fast enough as $x \rightarrow \pm \infty$ so that the normalization integral,



$$= E\psi(x)$$

Classical Picture

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Quantum Mechanical Realisation







$$x < 0 \quad \text{and} \quad x > L$$

$$\frac{d^2\psi(x)}{dx^2} + V(x)\psi(x) = E\psi(x)$$

$$\frac{h^2}{dx} \frac{d^2 \psi(x)}{dx^2} = E \psi(x)$$



 $\psi''(x) = -\frac{2mE}{\hbar^2}\psi(x) = -k^2\psi(x)$

where we have substituted the square of the wave number k, since

k

$\psi(x) = A \sin kx$



$$L^2 = \left(\frac{p}{\hbar}\right)^2 = \frac{2mE}{\hbar^2}$$



 $k_n = n \frac{\pi}{L}$

 $E_n = \frac{\hbar^2 k_n^2}{2m} =$

 $\int_{-\infty}^{+\infty} \psi^* \psi_n dx$

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$\psi(L) = A \sin kL = 0$

$$n = 1, 2, 3, \ldots$$

$$= n^2 \frac{\hbar^2 \pi^2}{2mL^2} = n^2 E_1$$

$$x = \int_0^L A_n^2 \sin^2\left(\frac{n\pi x}{L}\right) dx = 1$$

















The complete Wave function

$$e^{-i\omega t} = e^{-i(E_n)}$$

$$\Psi_n(x,t) = \sqrt{\frac{2}{L}} \mathrm{si}$$

$$\Psi_n(x,t) = \frac{1}{2i}\sqrt{\frac{2}{L}} \left[e^{i(k_n x - \omega_n t)} - e^{-i(k_n x + \omega_n t)} \right]$$

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 f_n/\hbar)t

 $\sin k_n x e^{-i\omega_n t}$

$$\sin k_n x = \frac{(e^{ik_n x} - e^{-ik_n x})}{2i}$$

An Electron in a Wire An electron moving in a thin metal wire is a reasonable approximation of a particle in a one-dimensional infinite well. The potential inside the wire is constant on average but rises sharply at each end. Suppose the electron is in a wire 1.0 cm long. (a) Compute the ground-state energy for the electron. (b) If the electron's energy is equal to the average kinetic energy of the molecules in a gas at T = 300 K, about 0.03 eV, what is the electron's quantum number n?

An Electron in an Atomic-Size Box (a) Find the energy in the ground state of an electron confined to a one-dimensional box of length L = 0.1 nm. (This box is roughly the size of an atom.) (b) Make an energy-level diagram and find the wavelengths of the photons emitted for all transitions beginning at state n = 3 or less and ending at a lower energy state.









